

DFT calculations of dichroic effect in transmission electron microscope

— ChiralTEM project —

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Outline

- 1 Introduction
 - Scattering cross-section
 - Mixed dynamic form-factor
- 2 Realization
 - Electronic Structure
 - Partial charges
 - BZ integration
 - Calculation of $S(q, q', E)$
- 3 Results
 - Face-centered cubic nickel
 - Body-centered cubic iron
 - Hexagonal close-packed cobalt
- 4 Conclusions and outlook

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Dynamical diffraction theory

- Electron in beam - $|\mathbf{k}\rangle$ -state
- Electron in crystal - superposition of Bloch waves
- Assuming crystal periodicity, boundary conditions lead to

$$\frac{d^2\sigma}{d\Omega dE} = \frac{4\gamma^2 k_{out}}{a_0^2 k_{in}} \sum_{\mathbf{J}} \sum_{\mathbf{J}'} T_{J_z, J'_z} A_{\mathbf{J}} A_{\mathbf{J}'}^* \times$$

$$\times \frac{1}{N} \sum_{\mathbf{u}} e^{i(\mathbf{Q}_{J_{xy}} - \mathbf{Q}_{J'_{xy}}) \cdot \mathbf{u}} \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$$

- Summation - a problem of geometry of experiment
- Mixed Dynamic Formfactor (MDFF) - a key quantity

MDFF properties

Symmetries

- $S(\mathbf{q}, \mathbf{q}', E) = \sum_{i,f} \langle i | e^{-i\mathbf{q}\cdot\mathbf{R}} | f \rangle \langle f | e^{i\mathbf{q}'\cdot\mathbf{R}} | i \rangle \delta(E + E_i - E_f)$
 - Clearly $S(\mathbf{q}, \mathbf{q}', E) = S^*(\mathbf{q}', \mathbf{q}, E)$
 - In non-magnetic systems: $S(\mathbf{q}, \mathbf{q}', E) = S(-\mathbf{q}', -\mathbf{q}, E)$
 - Centrosymmetric systems: $S(\mathbf{q}, \mathbf{q}', E) = S(-\mathbf{q}, -\mathbf{q}', E)$
-
- i.e., in non-magnetic centrosymmetric crystals
MDFF is real
 - But to observe dichroism, we need
nonzero imaginary part of MDFF

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Electronic Structure - WIEN2k program package

Self-consistency cycle

lapw0	Poisson equation ... $\Delta V = \rho$
lapw1	scal. rel. eigv. problem ... $HC^{(\nu)} = E^{(\nu)}SC^{(\nu)}$
lapwso	2 nd variational step ... $H_{SO}C_{SO}^{(\nu)} = E_{SO}^{(\nu)}C_{SO}^{(\nu)}$
lapw2	val. density ... $\rho(\mathbf{r}) = \sum_{\mathbf{k}\nu} \Theta(E_F - E_{\mathbf{k}\nu}) \psi_{\mathbf{k}\nu}(\mathbf{r}) ^2$
lcore	Dirac equation for core states
mixer	mixing old and new $\rho(\mathbf{r})$

Needed WIEN2k outputs

- core wavefunctions and energies
- bandstructure - Bloch functions $\psi_{\mathbf{k}\nu}(\mathbf{r})$ and energies $\epsilon_{\mathbf{k}\nu}$

Initial and final states in APW

Initial state

- Localized (core) state
- Nonzero only inside the atomic sphere

$$\langle \mathbf{r}, s | i \rangle \rightarrow \langle \mathbf{r}, s | jj_z \rangle = \sum_m C_{lm\frac{1}{2}s}^{jj_z} R_{js}(r) Y_l^m(\mathbf{r}/r)$$

Final state

- Itinerant Bloch state
- Defined in the whole unit cell; inside at. sphere we have

$$\langle \mathbf{r}, s | f \rangle \rightarrow \langle \mathbf{r}, s | \mathbf{k}\nu \rangle = \sum_{LM} D_{LMS}^{\mathbf{k}\nu} u_{LS}^{E_{\mathbf{k}\nu}}(r) Y_L^M(\mathbf{r}/r)$$

Mixed Dynamic Formfactor

$$\begin{aligned}
 S(\mathbf{q}, \mathbf{q}', E) &= \sum_{mm'} \sum_{LMS} \sum_{L'M'S'} \sum_{\lambda\mu} \sum_{\lambda'\mu'} 4\pi i^{\lambda-\lambda'} (2l+1) \sqrt{[\lambda, \lambda', L, L']} \\
 &\times Y_{\mu}^{\lambda}(\mathbf{q}/q)^* Y_{\mu'}^{\lambda'}(\mathbf{q}'/q') \langle j_{\lambda}(q) \rangle_{ELSj} \langle j_{\lambda'}(q') \rangle_{EL'S'j} \\
 &\times \begin{pmatrix} l & \lambda & L \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} l & \lambda' & L' \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} l & \lambda & L \\ -m & \mu & M \end{pmatrix} \begin{pmatrix} l & \lambda' & L' \\ -m' & \mu' & M' \end{pmatrix} \\
 &\times \sum_{j_z} (-1)^{m+m'} (2j+1) \begin{pmatrix} l & \frac{1}{2} & j \\ m & S & -j_z \end{pmatrix} \begin{pmatrix} l & \frac{1}{2} & j \\ m' & S' & -j_z \end{pmatrix} \\
 &\times \underbrace{\sum_{\nu\mathbf{k}} D_{LMS}(\nu\mathbf{k}) D_{L'M'S'}(\nu\mathbf{k})^* \delta(E + E_{n\ell\kappa} - E_{\nu\mathbf{k}})}_{\text{cross-DOS}}
 \end{aligned}$$

Partial charges - QTL code

Altered QTL code replaces LAPW2 -qtl

- Proper symmetrization of Bloch states \rightarrow full-BZ k-mesh no more necessary
- Spin-orbit interaction properly introduced
- Gives all 1056 independent components of $\langle lms | l'm's' \rangle_{\mathbf{k}\nu}$
- Improved accuracy of I/O operations

Running QTL

- Perl script `initqtl.pl` prepares input and definition files
- Executing `qtl qtl.def` leads to `case.qtl1` file containing cross-charges $D_{lms}(\mathbf{k}\nu)D_{l'm's'}^*(\mathbf{k}\nu)$

Integration over Brillouin zone - TETRA program

Changes introduced to original TETRA code

- Improved accuracy of internal and I/O operations
- Output of all cross-DOS components goes to a single file

Running TETRA

- Requires `case.inse1` file
- Perl script `initmdff.pl` prepares input and definition files for both TETRA and MDDFSO codes
- Executing `tetra tetra.def` leads to `case.crossdos` file - cross-DOS components

Cross-DOS components

Matrix 32x32 of complex numbers $\langle lms | l'm's' \rangle$ with $l_{max} = 3$:

$$\left(\begin{array}{cccc|cccc} \langle 00|00 \rangle & \langle 00|1\bar{1} \rangle & \dots & \langle 00|33 \rangle & \langle 00|00 \rangle & \langle 00|1\bar{1} \rangle & \dots & \langle 00|33 \rangle \\ \langle 1\bar{1}|00 \rangle & \langle 1\bar{1}|1\bar{1} \rangle & \dots & \langle 1\bar{1}|33 \rangle & \langle 1\bar{1}|00 \rangle & \langle 1\bar{1}|1\bar{1} \rangle & \dots & \langle 1\bar{1}|33 \rangle \\ \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \ddots & \vdots \\ \langle 33|00 \rangle & \langle 33|1\bar{1} \rangle & \dots & \langle 33|33 \rangle & \langle 33|00 \rangle & \langle 33|1\bar{1} \rangle & \dots & \langle 33|33 \rangle \\ \hline \langle 00|00 \rangle & \langle 00|1\bar{1} \rangle & \dots & \langle 00|33 \rangle & \langle 00|00 \rangle & \langle 00|1\bar{1} \rangle & \dots & \langle 00|33 \rangle \\ \langle 1\bar{1}|00 \rangle & \langle 1\bar{1}|1\bar{1} \rangle & \dots & \langle 1\bar{1}|33 \rangle & \langle 1\bar{1}|00 \rangle & \langle 1\bar{1}|1\bar{1} \rangle & \dots & \langle 1\bar{1}|33 \rangle \\ \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \ddots & \vdots \\ \langle 33|00 \rangle & \langle 33|1\bar{1} \rangle & \dots & \langle 33|33 \rangle & \langle 33|00 \rangle & \langle 33|1\bar{1} \rangle & \dots & \langle 33|33 \rangle \end{array} \right)$$

Spin blocks: $\left(\begin{array}{c|c} \uparrow\uparrow & \uparrow\downarrow \\ \hline \downarrow\uparrow & \downarrow\downarrow \end{array} \right)$

Calculation of $S(q,q',E)$ - MDFFSO program

Changes introduced to original $MDFE$ code

- Relativistic treatment of initial (core) states
- Extended to include spin-up/down cross-terms
- Calculation speed optimizations
- Improved accuracy of I/O operations (reading `case.crossdos` file)

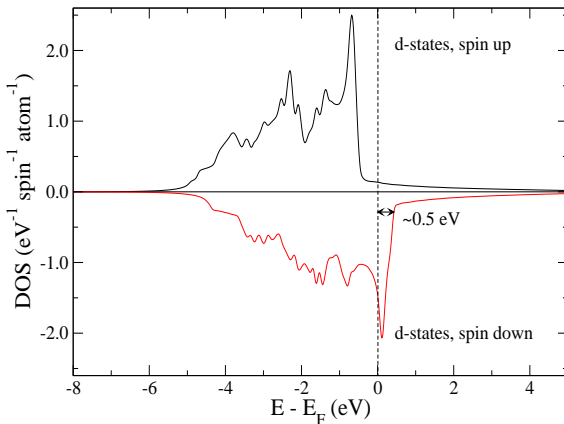
Running MDFFSO

- Requires `case.inse1` file
- Executing `mdffso mdffso.def` leads to `case.mdffso` file - ready for summation with `ilX` codes

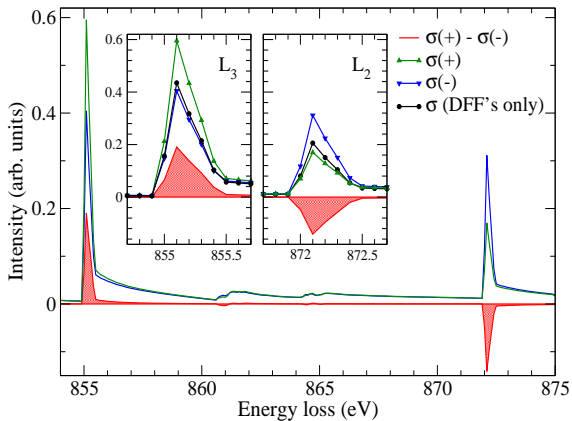
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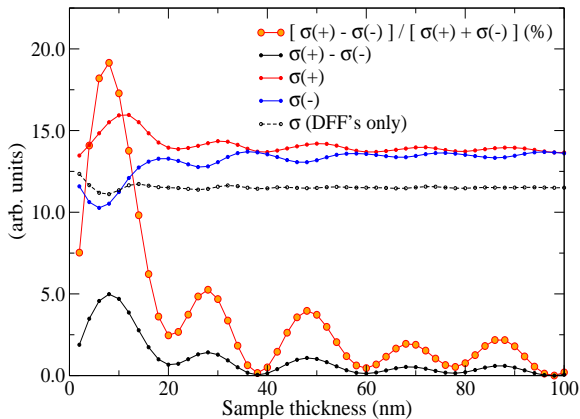
fccNi: density of states - 3d projections



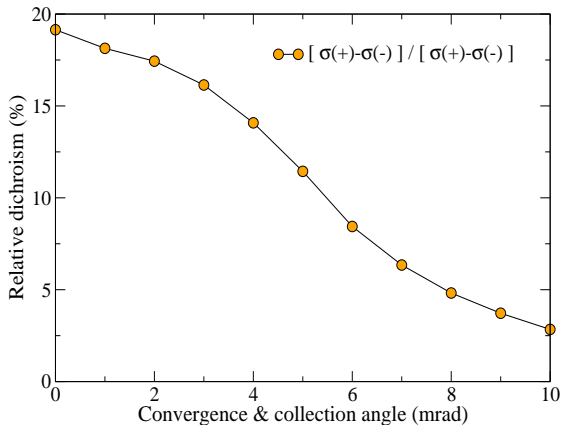
fccNi: dichroism



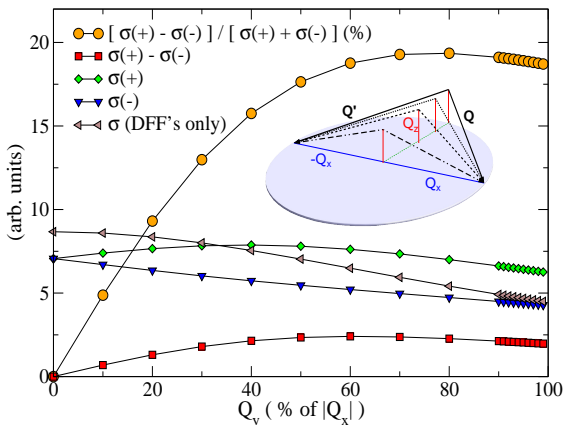
fccNi: thickness profile



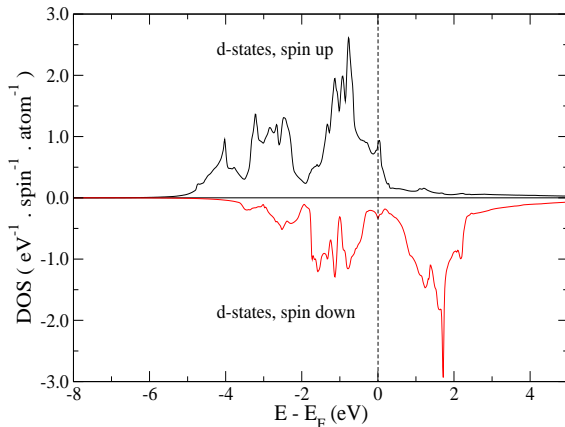
fccNi: convergence & collection angle profile



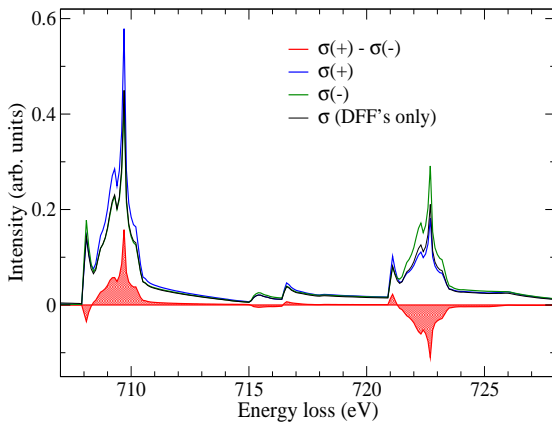
fccNi: detector placement sensitivity



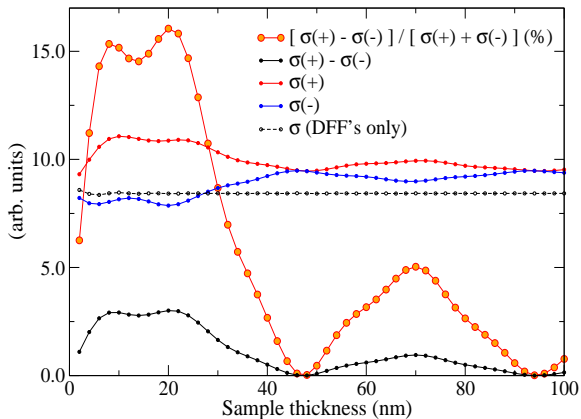
bccFe: density of states - 3d projections



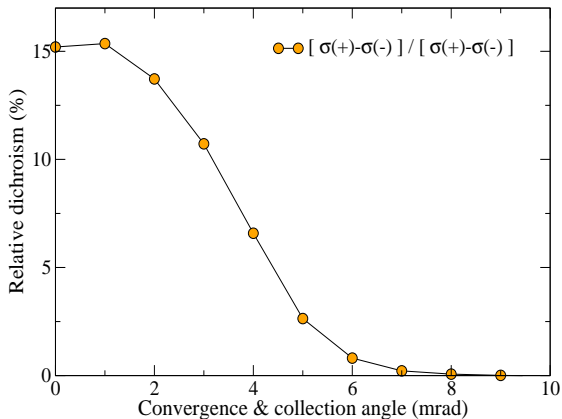
bccFe: dichroism



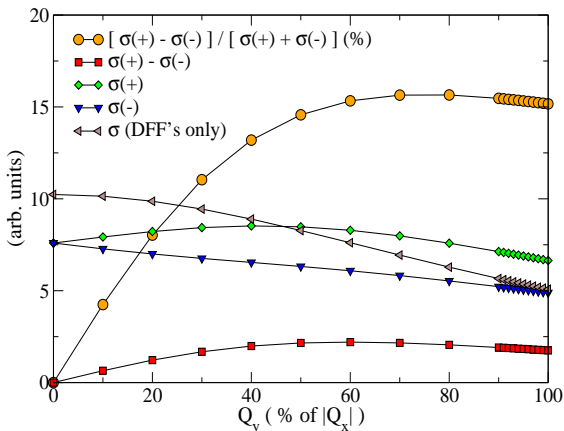
bccFe: thickness profile



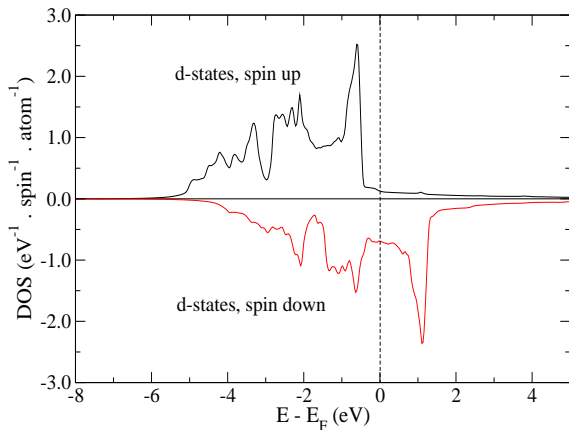
bccFe: convergence & collection angle profile



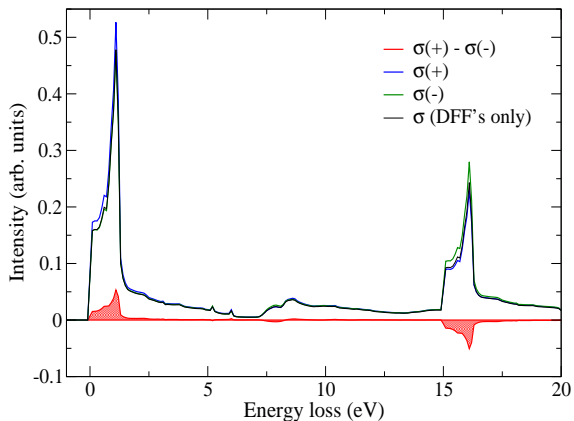
bccFe: detector placement sensitivity



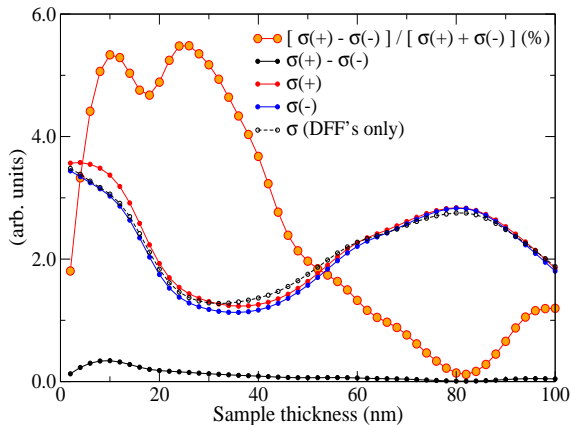
hcpCo: density of states - 3d projections



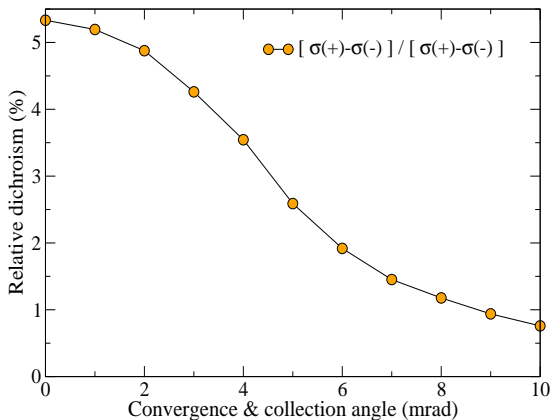
hcpCo: dichroism



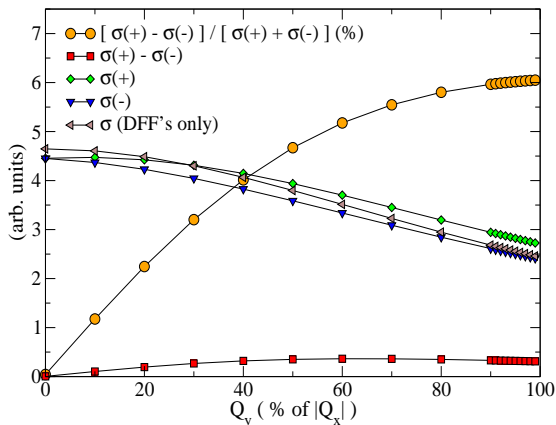
hcpCo: thickness profile



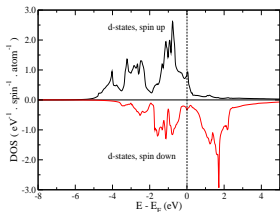
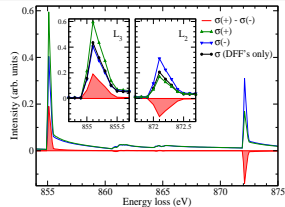
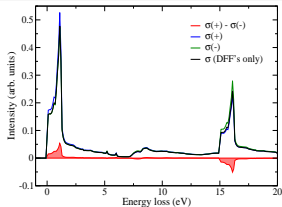
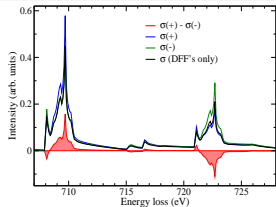
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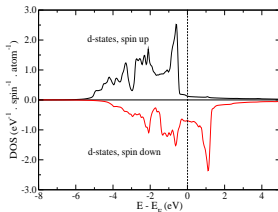
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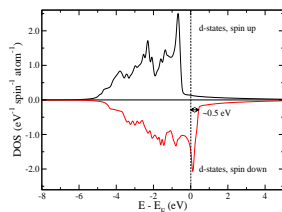
Why so small dichroism in Co?



bcc Fe



hcp Co



fcc Ni

Why so small dichroism in Co?

Effect of summation over atoms

$$\frac{d^2\sigma}{d\Omega dE} = \frac{4\gamma^2}{a_0^2} \frac{k_{out}}{k_{in}} \sum_{\mathbf{J}} \sum_{\mathbf{J}'} T_{J_z, J'_z} A_{\mathbf{J}} A_{\mathbf{J}'}^* \frac{1}{N} \sum_{\mathbf{u}} e^{i(\mathbf{Q}_{J_{xy}} - \mathbf{Q}_{J'_{xy}}) \cdot \mathbf{u}} \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$$

- Ni & Fe: $N = 1$, $\mathbf{u} = 0 \rightarrow \frac{1}{N} \sum(\dots) = \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$

- Co: $N = 2$, $\mathbf{u}_{1,2} = \{\frac{1}{3}, \frac{2}{3}, \frac{1}{4}\}, \{\frac{2}{3}, \frac{1}{3}, \frac{3}{4}\}$

DFF: $\mathbf{q} = \mathbf{q}' \rightarrow \frac{1}{N} \sum(\dots) = 1 \times \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$

MDFF: $\mathbf{q} \perp \mathbf{q}'$, $\mathbf{q} - \mathbf{q}' = \mathbf{G}$, $\mathbf{G} = (2, 0, 0)$

$\exp(i\mathbf{u}_1 \cdot \mathbf{G}) = (-\frac{1}{2}, \frac{\sqrt{3}}{2})$, $\exp(i\mathbf{u}_2 \cdot \mathbf{G}) = (-\frac{1}{2}, -\frac{\sqrt{3}}{2})$

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MDFF: $\mathbf{q} \perp \mathbf{q}'$, $\mathbf{q} - \mathbf{q}' = \mathbf{G}$, $\mathbf{G} = (2, 0, 0)$

$\exp(i\mathbf{u}_1 \cdot \mathbf{G}) = (-\frac{1}{2}, \frac{\sqrt{3}}{2})$, $\exp(i\mathbf{u}_2 \cdot \mathbf{G}) = (-\frac{1}{2}, -\frac{\sqrt{3}}{2})$

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Why so small dichroism in Co?

Effect of summation over atoms

$$\frac{d^2\sigma}{d\Omega dE} = \frac{4\gamma^2}{a_0^2} \frac{k_{out}}{k_{in}} \sum_{\mathbf{J}} \sum_{\mathbf{J}'} T_{J_z, J'_z} A_{\mathbf{J}} A_{\mathbf{J}'}^* \frac{1}{N} \sum_{\mathbf{u}} e^{i(\mathbf{Q}_{J_{xy}} - \mathbf{Q}_{J'_{xy}}) \cdot \mathbf{u}} \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$$

- Ni & Fe: $N = 1$, $\mathbf{u} = 0 \rightarrow \frac{1}{N} \sum(\dots) = \frac{S_{\mathbf{u}}(\mathbf{Q}_{J_{xy}}, \mathbf{Q}_{J'_{xy}}, E)}{Q_{J_{xy}}^2 Q_{J'_{xy}}^2}$

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Can we do better?

Yes, let's choose a better \mathbf{G}

- let $\mathbf{G} = (G_x, G_y, 0)$
- requiring $\mathbf{G} \cdot \mathbf{u}_1 = 2\pi n_1$ and $\mathbf{G} \cdot \mathbf{u}_2 = 2\pi n_2$
we get (in (hkl) notation) $\mathbf{G} = (2n_2 - n_1, 2n_1 - n_2, 0)$
- e.g., let $n_1 = 1$ and $n_2 = 1$, then $\mathbf{G} = (110)$.

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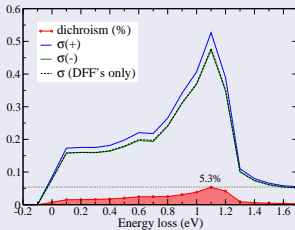
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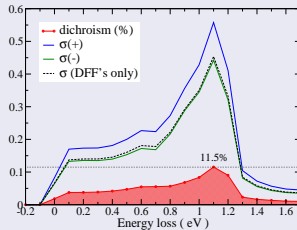
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$\mathbf{G} = (200)$



$\mathbf{G} = (110)$

Outline

- 1 Introduction
 - Scattering cross-section
 - Mixed dynamic form-factor
- 2 Realization
 - Electronic Structure
 - Partial charges
 - BZ integration
 - Calculation of $S(q, q', E)$
- 3 Results
 - Face-centered cubic nickel
 - Body-centered cubic iron
 - Hexagonal close-packed cobalt
- 4 **Conclusions and outlook**

Conclusions and outlook

Conclusions

- A coherent set of programs for calculation of MDFF was prepared
- A dichroic effect on bccFe, hcpCo and fccNi was calculated
- Calculations indicate that dichroic effect is:
 - weakly sensitive to detector misplacements of up to 30% q_y
 - rather sensitive to sample thickness
 - sensitive to collection/convergence angle
- There is an important dependence of dichroism strength on exp. geometry - on the systematic row indices